

Variable separation method.

Examples

- Using variable separation, solve Laplace's equation $\nabla^2 u = 0$ in two dimensions using plane polar coordinates,

$$\frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u}{\partial \varphi^2} = 0,$$

and show that a solution of this equation can be constructed as

$$u(r, \varphi) = C \ln r + D + \sum_{n=1}^{\infty} (A_n \cos n\varphi + B_n \sin n\varphi)(E_n r^n + F_n r^{-n}),$$

where C, D, A_n, B_n, E_n and F_n are arbitrary constants.

- Prove that

$$\int_0^l \sin \frac{n\pi x}{l} \sin \frac{m\pi x}{l} dx = \begin{cases} 0 & n = m = 0 \\ \frac{1}{2} l \delta_{nm} & \text{otherwise} \end{cases} \quad (1)$$

$$\int_{-l}^l \sin \frac{n\pi x}{l} \cos \frac{m\pi x}{l} dx = 0 \quad (\text{note the limits}) \quad (2)$$

$$\int_0^l \cos \frac{n\pi x}{l} \cos \frac{m\pi x}{l} dx = \begin{cases} l & n = m = 0 \\ \frac{1}{2} l \delta_{nm} & \text{otherwise,} \end{cases} \quad (3)$$

for integer $n, m \geq 0$, where $\delta_{nm} = 1$ for $n = m$, 0 for $n \neq m$, is the *Kronecker delta* symbol.

Homework problems

- Using the form $u(x, t) = v(x)q(t)$, solve the one-dimensional wave equation,

$$\frac{\partial^2 u}{\partial t^2} - c^2 \frac{\partial^2 u}{\partial x^2} = 0,$$

for the string of length l ($0 \leq x \leq l$) with boundary conditions $u(0, t) = 0, u_x(l, t) = 0$ (i.e., fixed end at $x = 0$ and "free" end at $x = l$).

Hence, show that the string can execute harmonic vibrations described by

$$u(x, t) = A \sin \left[\pi \left(n + \frac{1}{2} \right) x / l \right] \cos(\omega_n t + \phi),$$

with frequencies $\omega_n = \pi c \left(n + \frac{1}{2} \right) / l, n = 0, 1, \dots$, and arbitrary amplitude A and phase ϕ .

- A rope of length l and linear mass density ρ hangs freely along the x axis under gravity (acceleration g). The bottom end of the string lies at $x = 0$ and the top at $x = l$.

- Using the approach used for the string, show that the displacement $u(x, t)$ of the rope satisfies the equation

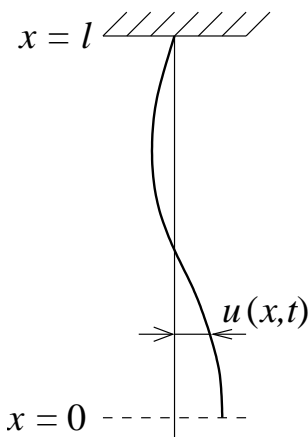
$$\frac{\partial^2 u}{\partial t^2} - g \frac{\partial}{\partial x} \left(x \frac{\partial u}{\partial x} \right) = 0. \quad (4)$$

[Hint: at point x the tension force in the rope is $T = g\rho x$.]

- Seeking solution of Eq. (4) in the form $u(x, t) = v(x)q(t)$, find $q(t)$ and show that $v(x)$ satisfies the equation

$$x \frac{d^2 v}{dx^2} + \frac{dv}{dx} + \frac{\omega^2}{g} v = 0, \quad (5)$$

where $-\omega^2$ is the separation constant.



- (c) Introduce a new independent variable $\xi = \alpha\sqrt{x}$, i.e., $x = \xi^2/\alpha^2$, where α is a constant, and show that Eq. (5) takes the form

$$\frac{d^2v}{d\xi^2} + \frac{1}{\xi} \frac{dv}{d\xi} + v = 0, \quad (6)$$

if one chooses $\alpha = 2\omega/\sqrt{g}$.¹

- (d) Equation (6) is the Bessel equation for $m = 0$, whose regular solution is $J_0(\xi)$. Hence, show that the solutions $v(x)$ of Eq. (5) such that $v(0)$ is finite and $v(l) = 0$, are

$$v(x) = AJ_0\left(z_{0,n}\sqrt{\frac{x}{l}}\right), \quad n = 1, 2, \dots, \quad (7)$$

where A is an arbitrary constant, $z_{0,n}$ is the n th root of $J_0(z)$, and $\omega \equiv \omega_n = \frac{z_{0,n}}{2}\sqrt{\frac{g}{l}}$.

- (e) Combining the results from (a)–(d), show that the hanging rope executing harmonic motion with frequency ω_n , is described by

$$u(x, t) = AJ_0\left(z_{0,n}\sqrt{\frac{x}{l}}\right) \cos(\omega_n t + \phi),$$

where ϕ is an arbitrary initial phase.

3. Consider the one-dimensional heat equation for $0 \leq x \leq l$ (rod of length l),

$$u_t - Ku_{xx} = 0. \quad (8)$$

- (a) Show that when the rod is in thermal equilibrium (i.e., the temperature does not change with time, $\partial u/\partial t = 0$), the time-independent (or *stationary*) solution of Eq. (8), $u_s(x)$, which satisfies the boundary conditions $u_s(0) = T_1$, $u_s(l) = T_2$, is

$$u_s(x) = T_1 + (T_2 - T_1)x/l. \quad (9)$$

- (b) Show that if $u_0(x, t)$ is a solution of Eq. (8) with $u_0(0, t) = u_0(l, t) = 0$, then $u(x, t) = u_0(x, t) + u_s(x)$ satisfies Eq. (8) with boundary conditions $u(0, t) = T_1$, $u(l, t) = T_2$.

- (c) Using

$$u_0(x, t) = \sum_{n=1}^{\infty} B_n \sin \frac{n\pi x}{l} e^{-(n^2\pi^2/l^2)Kt}, \quad (10)$$

show that the solution which satisfies $u(0, t) = T_1$, $u(l, t) = T_2$ and the *initial* condition $u(x, 0) = f(x)$, is

$$u(x, t) = T_1 + (T_2 - T_1)x/l + \sum_{n=1}^{\infty} B_n \sin \frac{n\pi x}{l} e^{-(n^2\pi^2/l^2)Kt}, \quad (11)$$

where $B_n = \frac{2}{l} \int_0^l \sin \frac{n\pi x}{l} [f(x) - T_1 - (T_2 - T_1)x/l] dx$.

4. Using the method of separation of variables, solve the two-dimensional wave equation in Cartesian coordinates for a rectangular membrane ($0 \leq x \leq a$, $0 \leq y \leq b$) with fixed edges, $u(0, y, t) = u(a, y, t) = u(x, 0, t) = u(x, b, t)$, and show that the membrane executes harmonic motion with frequencies $\omega_{nm} = \pi c(n^2/a^2 + m^2/b^2)^{1/2}$, where $n, m = 0, 1, 2, \dots$, and described by $u(x, y, t) = A \sin(n\pi x/a) \sin(m\pi y/b) \cos(\omega_{nm}t + \phi)$.

¹Hint: use chain rule to transform the derivatives, $\frac{dv}{dx} = \frac{dv}{d\xi} \frac{d\xi}{dx}$, $\frac{d^2v}{dx^2} = \frac{d^2v}{d\xi^2} \left(\frac{d\xi}{dx}\right)^2 + \frac{dv}{d\xi} \frac{d^2\xi}{dx^2}$.