

Examples

①  $u_{xx} + u_{yy} = 0$  for  $u(x,y)$ ,  $-\infty < x < \infty$ ,  
 $0 \leq y < \infty$ .

Boundary condition:  $u(x,0) = f(x)$ .

(a)  $U(p,y) = \mathcal{F}[u] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ipx} u(x,y) dx$

1)  $\mathcal{F}[u_{yy}] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ipx} \frac{\partial^2 u(x,y)}{\partial y^2} dx$   
 $= \frac{\partial^2}{\partial y^2} \left( \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ipx} u(x,y) dx \right)$

The order of integration over  $x$  and differentiation with respect to  $y$  can be interchanged.

$\Rightarrow \mathcal{F}[u_{yy}] = \frac{\partial^2 U(p,y)}{\partial y^2}$

2)  $\mathcal{F}[u_{xx}] = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ipx} \frac{\partial^2 u}{\partial x^2} dx$

Here we use integration by parts:  
 $\int f g' dx = f g - \int f' g dx$

$= \frac{1}{\sqrt{2\pi}} e^{ipx} \frac{\partial u}{\partial x} \Big|_{-\infty}^{+\infty} - \frac{ip}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ipx} \frac{\partial u}{\partial x} dx$

since  $\frac{\partial u}{\partial x} \xrightarrow{x \rightarrow \pm\infty} 0$

$= -\frac{ip}{\sqrt{2\pi}} e^{ipx} u \Big|_{-\infty}^{+\infty} + \frac{(ip)^2}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ipx} u(x,y) dx$

since  $u(x,y) \xrightarrow{x \rightarrow \pm\infty} 0$

$= -\frac{p^2}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ipx} u(x,y) dx$

Therefore,  $\mathcal{F}[u_{xx}] = -p^2 U(p,y)$

(b) Applying the Fourier transform to the Laplace equation:  $\mathcal{F}[u_{xx}] + \mathcal{F}[u_{yy}] = 0$ , gives: (2)

$$\frac{\partial^2 U}{\partial y^2} - p^2 U = 0 \quad (*)$$

This is an ordinary differential equation with respect to  $y$ ,  $p$  being constant (a parameter).

The general solution of (\*) is

$$U(p, y) = A(p) e^{-py} + B(p) e^{py}$$

The solution we seek should not diverge at  $y \rightarrow \infty$ . Hence, for  $p > 0$  we must choose  $e^{-py}$ , while for  $p < 0$  we should choose  $e^{py}$ .

This can be done by writing

$$U(p, y) = A(p) e^{-|p|y} \quad (**)$$

Let us now find  $A(p)$ . We know from the boundary condition that  $u(x, 0) = f(x)$ .

Taking the Fourier transform of this equation, we have:

$$U(p, 0) = F(p), \text{ where } F(p) = \mathcal{F}[f].$$

By setting  $y = 0$  in (\*\*), we obtain:

$$A(p) = F(p),$$

hence 
$$U(p, y) = F(p) e^{-|p|y}.$$

To find  $u(x, y)$  we must perform the inverse transform of  $U(p, y)$ .

$$u(x,y) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-ipx} U(p,y) dp$$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{-ipx} F(p) e^{-|p|y} dp$$

$$= \left(\frac{1}{\sqrt{2\pi}}\right)^2 \int_{-\infty}^{+\infty} e^{-ipx} e^{-|p|y} dp \int_{-\infty}^{+\infty} e^{ip\xi} f(\xi) d\xi$$

$F(p) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{+\infty} e^{ip\xi} f(\xi) d\xi$   
 where we use  $\xi$  to avoid confusion when this is substituted.

Changing the order of integration,

$$= \frac{1}{2\pi} \int_{-\infty}^{+\infty} f(\xi) d\xi \int_{-\infty}^{+\infty} e^{ip(\xi-x)-|p|y} dp$$

To expand the modulus,  $|p|$ , we split the integral into two, with  $p \geq 0$  and  $p \leq 0$ .

$$= \int_{-\infty}^0 e^{ip(\xi-x)+py} dp + \int_0^{+\infty} e^{ip(\xi-x)-py} dp$$

$|p| = p$  for  $p \geq 0$   
 $|p| = -p$  for  $p \leq 0$

$$= \int_{-\infty}^0 e^{p(y+i(\xi-x))} dp + \int_0^{+\infty} e^{-p(y-ip(\xi-x))} dp$$

$$= \frac{e^{p(y+i(\xi-x))}}{y+i(\xi-x)} \Big|_{-\infty}^0 - \frac{e^{-p(y-ip(\xi-x))}}{y-i(\xi-x)} \Big|_0^{+\infty}$$

only upper limit contributes here

only lower limit contributes here

$$= \frac{1}{y+i(\xi-x)} + \frac{1}{y-i(\xi-x)} = \frac{y-i(\xi-x) + y+i(\xi-x)}{(y+i(\xi-x))(y-i(\xi-x))}$$

$$= \frac{2y}{y^2 + (\xi-x)^2}$$

Substituting this into the integral over  $\xi$ :

$$u(x,y) = \frac{y}{\pi} \int_{-\infty}^{+\infty} \frac{f(\xi) d\xi}{y^2 + (\xi-x)^2}$$

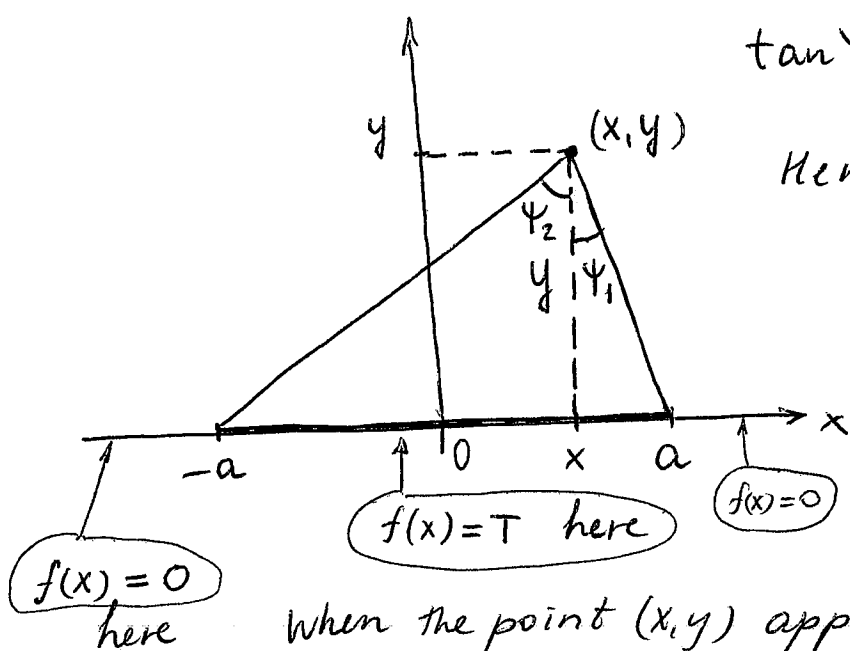
(c) Let us calculate  $u(x,y)$  for the special case  $f(x) = T$  for  $-a \leq x \leq a$ , and 0 otherwise. Substituting this in the last equation,

$$\begin{aligned}
 u(x,y) &= \frac{y}{\pi} \int_{-a}^a \frac{T d\xi}{y^2 + (\xi - x)^2} \\
 &= \frac{yT}{\pi} \int_{-a}^a \frac{d\xi}{y^2 + (\xi - x)^2} \\
 &= \frac{yT}{\pi} \frac{1}{y^2} \int_{-a}^a \frac{d\xi}{1 + \left(\frac{\xi - x}{y}\right)^2} \quad \left. \begin{array}{l} \text{We know} \\ \text{a table integral} \\ \int \frac{dx}{1+x^2} = \arctan x \end{array} \right\} \\
 &= \frac{yT}{\pi} \frac{y}{y^2} \int_{-a}^a \frac{d\left(\frac{\xi - x}{y}\right)}{1 + \left(\frac{\xi - x}{y}\right)^2} \\
 &= \frac{T}{\pi} \left[ \arctan \frac{\xi - x}{y} \right]_{-a}^a = \frac{T}{\pi} \left[ \arctan \frac{a-x}{y} - \arctan \frac{-a-x}{y} \right]
 \end{aligned}$$

$\arctan$  is an odd function, hence we can write:

$$u(x,y) = \frac{T}{\pi} \left[ \arctan \frac{a-x}{y} + \arctan \frac{a+x}{y} \right]$$

This result has an interesting geometric interpretation.



$$\tan \Psi_1 = \frac{a-x}{y}, \quad \tan \Psi_2 = \frac{a+x}{y}$$

Hence

$$u(x,y) = \frac{T}{\pi} [\Psi_1 + \Psi_2],$$

and  $\Psi_1 + \Psi_2$  is the angle subtended by the  $(-a, a)$  segment of the  $x$  axis.

When the point  $(x,y)$  approaches the  $x$  axis within  $-a < x < a$ ,  $\Psi_1 + \Psi_2 \rightarrow \pi$  and  $u(x,y) \rightarrow T$ , as required.

(2) Heat equation:  $u_t - u_{xx} = 0$  for  $u(x,t)$ , (5)  
 $0 \leq x < \infty$ , with initial condition  $u(x,0) = 0$   
 and boundary condition  $u_x(0,t) = -5$ .

[ Recall:  $j = -\alpha \frac{\partial u}{\partial x}$  is the heat flux, so that  
 $\frac{\partial u}{\partial x} \Big|_{x=0} = -5$  means  $j = +\alpha 5$  - constant heat  
 flux into the rod. ]

Using the cosine Fourier transform:

$$U_c(p,t) = \mathcal{F}_c [u] = \sqrt{\frac{2}{\pi}} \int_0^{\infty} u(x,t) \cos px \, dx$$

Fourier transform of  $u_t$ :

$$\mathcal{F} [u_t] = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\partial u}{\partial t} \cos px \, dx = \frac{\partial}{\partial t} \left( \sqrt{\frac{2}{\pi}} \int_0^{\infty} u(x,t) \cos px \, dx \right)$$

$$\Rightarrow \mathcal{F} [u_t] = \frac{\partial U_c(p,t)}{\partial t}$$

Fourier transform of  $u_{xx}$ :

$$\mathcal{F} [u_{xx}] = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\partial^2 u}{\partial x^2} \cos px \, dx$$

} Integrating  
by parts

$$= \sqrt{\frac{2}{\pi}} \frac{\partial u}{\partial x} \cos px \Big|_0^{\infty} + \sqrt{\frac{2}{\pi}} p \int_0^{\infty} \frac{\partial u}{\partial x} \sin px \, dx$$

$$= -\sqrt{\frac{2}{\pi}} u_x(0,t) + \underbrace{\sqrt{\frac{2}{\pi}} p u(x,t) \sin px \Big|_0^{\infty}}_{\text{This is zero both at } x=\infty \text{ and } x=0} - \sqrt{\frac{2}{\pi}} p^2 \int_0^{\infty} u(x,t) \cos px \, dx$$

only the lower limit  
 $x=0$  contributes here;

we assume  $\frac{\partial u}{\partial x} \rightarrow 0$   
 $x \rightarrow \infty$

and  $u \rightarrow 0$   
 $x \rightarrow \infty$

$$\text{Therefore: } \mathcal{F} [u_{xx}] = -\sqrt{\frac{2}{\pi}} u_x(0,t) - p^2 U_c(p,t)$$

Let us apply the cosine Fourier transform to (6) the heat equation:  $\mathcal{F}[u_t] - \mathcal{F}[u_{xx}] = 0$ ,

which gives:  $\frac{\partial U_c(p,t)}{\partial t} + \sqrt{\frac{2}{\pi}} u_x(0,t) + p^2 U_c(p,t) = 0$ ,

or  $\frac{\partial U_c(p,t)}{\partial t} - \sqrt{\frac{2}{\pi}} \sigma + p^2 U_c(p,t) = 0$ .

$$\frac{\partial U_c}{\partial t} + p^2 U_c = \sqrt{\frac{2}{\pi}} \sigma \quad (*)$$

This is a first-order linear ordinary differential equation. The homogeneous equation is

$$\frac{\partial U_c}{\partial t} + p^2 U_c = 0,$$

has the solution  $U_c = C(p) e^{-p^2 t}$ .

To solve the inhomogeneous equation, replace the "constant"  $C(p)$  with a function of  $t$ ,  $C(p,t)$ .

$$U_c = C(p,t) e^{-p^2 t}.$$

Substituting into (\*):

$$\frac{\partial C}{\partial t} e^{-p^2 t} + C(p,t) (-p^2) e^{-p^2 t} + p^2 C(p,t) e^{-p^2 t} = \sqrt{\frac{2}{\pi}} \sigma$$

$$\frac{\partial C}{\partial t} = \sqrt{\frac{2}{\pi}} \sigma e^{p^2 t}$$

Integrating over  $t$ :

$$C(p,t) = \sqrt{\frac{2}{\pi}} \frac{\sigma}{p^2} e^{p^2 t} + C_1(p)$$

↑ arbitrary constant  
(function of  $p$  only)

$$\Rightarrow U_c(p,t) = \left( \sqrt{\frac{2}{\pi}} \frac{\sigma}{p^2} e^{p^2 t} + C_1(p) \right) e^{-p^2 t}$$

$$U_c(p,t) = C_1(p) e^{-p^2 t} + \sqrt{\frac{2}{\pi}} \frac{\sigma}{p^2} \quad (**)$$

Using the initial condition,  $u(x,0) = 0$ , (7)  
we obtain:  $U_c(p,0) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} u(x,0) \cos px dx = 0$ .

Setting  $t = 0$  in equation (\*\*) (page 6):

$$C_1(p) + \sqrt{\frac{2}{\pi}} \frac{\sigma}{p^2} = 0 \Rightarrow C_1(p) = -\sqrt{\frac{2}{\pi}} \frac{\sigma}{p^2}.$$

Hence:

$$U_c(p,t) = \sqrt{\frac{2}{\pi}} \sigma \frac{1 - e^{-p^2 t}}{p^2}.$$

From this,  $u(x,t)$  is obtained by performing the inverse cosine Fourier transform:

$$u(x,t) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} U_c(p,t) \cos px dp,$$

i.e.,

$$u(x,t) = \frac{2\sigma}{\pi} \int_0^{\infty} \frac{1 - e^{-p^2 t}}{p^2} \cos px dp.$$

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This integral can, in fact, be expressed in terms of the error function,  $\text{erf}(z) = \frac{2}{\sqrt{\pi}} \int_0^z e^{-s^2} ds$ ,

as shown in Problem 355 of the book by N.N. Lebedev, I.P. Skal'skaya and Ya.S. Uflyand, *Problems in Mathematical Physics* (Pergamon Press, Oxford, 1966):

$$u(x,t) = \sigma \left\{ \frac{2\sqrt{t}}{\sqrt{\pi}} e^{-\frac{x^2}{4t}} - x \left[ 1 - \text{erf}\left(\frac{x}{2\sqrt{t}}\right) \right] \right\}.$$

In fact, this is not too difficult to demonstrate (see next page if you are curious!).

Let's differentiate  $u(x,t)$  with respect to  $t$ : (8)

$$\begin{aligned} \frac{\partial u(x,t)}{\partial t} &= \frac{2\sigma}{\pi} \int_0^{\infty} e^{-p^2 t} \cos px \, dp && \text{this is an even function,} \\ &= \frac{2\sigma}{\pi} \frac{1}{2} \int_{-\infty}^{\infty} e^{-p^2 t} (e^{ipx} + e^{-ipx}) \, dp && \text{change } \int_0^{\infty} \text{ to } \int_{-\infty}^{\infty} \\ &= \frac{\sigma}{2\pi} \left[ \int_{-\infty}^{\infty} e^{-tp^2 + ipx} \, dp + \int_{-\infty}^{\infty} e^{-tp^2 - ipx} \, dp \right] \end{aligned}$$

Using  $\int_{-\infty}^{\infty} e^{-\alpha x^2 - \beta x} \, dx = \sqrt{\frac{\pi}{\alpha}} e^{-\frac{\beta^2}{4\alpha}}$

$$= \frac{\sigma}{2\pi} \left( \sqrt{\frac{\pi}{t}} e^{\frac{(ix)^2}{4t}} + \sqrt{\frac{\pi}{t}} e^{\frac{(-ix)^2}{4t}} \right) = \frac{\sigma}{\sqrt{\pi}} \frac{1}{\sqrt{t}} e^{-\frac{x^2}{4t}}$$

To obtain  $u(x,t)$  we need to integrate this with respect to  $t$  from 0 to  $t$  ( $u(x,0) = 0$ ).

$$\Rightarrow u(x,t) = \frac{\sigma}{\sqrt{\pi}} \int_0^t \frac{1}{\sqrt{t}} e^{-\frac{x^2}{4t}} \, dt$$

Let us introduce the new variable  $s^2 = \frac{x^2}{4t}$ ,  $s = \frac{x}{2\sqrt{t}}$ .

$$ds = -\frac{x}{2 \cdot 2 t^{3/2}} dt \quad \text{or} \quad t = \frac{x^2}{4s^2}, \quad dt = -\frac{x^2}{2s^3} ds,$$

so that:  $u(x,t) = \frac{\sigma}{\sqrt{\pi}} \int \frac{2s}{x} e^{-s^2} \left(-\frac{x^2}{2s^3}\right) ds$

$$= \frac{\sigma}{\sqrt{\pi}} \frac{x^2}{x} \int_{x/2\sqrt{t}}^{\infty} e^{-s^2} \frac{ds}{s^2} = \frac{\sigma x}{\sqrt{\pi}} \int_{x/2\sqrt{t}}^{\infty} e^{-s^2} d\left(-\frac{1}{s}\right) \quad \left. \begin{array}{l} \text{Integrating by} \\ \text{parts} \end{array} \right\}$$

$$= \frac{\sigma x}{\sqrt{\pi}} \left[ e^{-s^2} \left(-\frac{1}{s}\right) \Big|_{x/2\sqrt{t}}^{\infty} + \int_{x/2\sqrt{t}}^{\infty} \frac{1}{s} e^{-s^2} (-2s) ds \right] \quad \left. \begin{array}{l} \int_0^{\infty} e^{-s^2} ds = \frac{\sqrt{\pi}}{2} \\ \frac{2}{\sqrt{\pi}} \int_0^{\infty} e^{-s^2} ds = 1 \end{array} \right\}$$

$$= \frac{\sigma x}{\sqrt{\pi}} \left[ +e^{-\frac{x^2}{4t}} \frac{2\sqrt{t}}{x} - \frac{2\sqrt{t}}{\sqrt{\pi}} \left( \int_0^{\infty} - \int_0^{x/2\sqrt{t}} \right) e^{-s^2} ds \right]$$

$$= \sigma \left[ \frac{2\sqrt{t}}{\sqrt{\pi}} e^{-\frac{x^2}{4t}} - x \left( 1 - \operatorname{erf} \left( \frac{x}{2\sqrt{t}} \right) \right) \right], \quad \text{as in the book quoted.}$$